

Spin^T structure and Dirac operator on Riemannian manifolds

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Abstract

In this paper, we describe the group $\text{Spin}^T(n)$ and give some properties of this group. We construct Spin^T spinor bundle \mathbb{S} by means of the spinor representation of the group $\text{Spin}^T(n)$ and define covariant derivative operator and Dirac operator on \mathbb{S} . Finally, Schrödinger-Lichnerowicz-type formula is derived by using these operators.

Key Words Spinor bundle, the group $\text{Spin}^T(n)$, Dirac operator, Schrödinger-Lichnerowicz-type formula.

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1 Introduction

Spin and Spin^c structures is effective tool to study the geometry and topology of manifolds, especially in dimension four. Spin and Spin^c manifolds have been studied extensively in [2, 3, 4, 5]. For any compact Lie group G the Spin^G structure have been studied in [1]. However, the spinor representation is replaced by a hyperkahler manifold, also called target manifold. In this paper, we define the Lie group $\text{Spin}^T(n)$ as a quotient group by taking $G = S^1 \times S^1$. The groups $\text{Spin}(n)$ and $\text{Spin}^c(n)$ are the subset of $\text{Spin}^T(n)$. We define Spin^T structure on any Riemannian manifold. The spinor representation of $\text{Spin}^T(n)$ is defined by the help of the spinor representation of $\text{Spin}(n)$. By using the spinor representation of $\text{Spin}^T(n)$ we construct the Spin^T spinor bundle \mathbb{S} . Finally, we give Schrödinger-Lichnerowicz-type formula by using covariant derivative operator and Dirac operator on \mathbb{S} .

This paper is organized as follows. We begin with a section introducing the group $\text{Spin}^T(n)$. In the following section, we define Spin^T structure on any Riemannian manifold. The final section is dedicated to the construction of

the spinor bundle \mathbb{S} , the study of the Dirac operator associated to Levi-Civita connection ∇ and Schrödinger-Lichnerowicz-type formula.

2 The group $\text{Spin}^T(n)$

Definition 1 *The Spin^T group is defined as*

$$\text{Spin}^T(n) := (\text{Spin}(n) \times S^1 \times S^1) / \{\pm 1\}.$$

The elements of $\text{Spin}^T(n)$ are thus classes $[g, z_1, z_2]$ of pairs $(g, z_1, z_2) \in \text{Spin}(n) \times S^1 \times S^1$ under the equivalence relation

$$(g, z_1, z_2) \sim (-g, -z_1, -z_2).$$

We can define the following homomorphisms:

- a. The map $\lambda^T : \text{Spin}^T(n) \longrightarrow SO(n)$ is given by $\lambda^T([g, z_1, z_2]) = \lambda(g)$ where the map $\lambda : \text{Spin}(n) \rightarrow SO(n)$ is the two-fold covering given by $\lambda(g)(v) = gv g^{-1}$.
- b. $i : \text{Spin}(n) \longrightarrow \text{Spin}^T(n)$ is the natural inclusion map $i(g) = [g, 1, 1]$.
- c. $j : S^1 \times S^1 \longrightarrow \text{Spin}^T(n)$ is the inclusion map $j(z_1, z_2) = [1, z_1, z_2]$.
- d. $l : \text{Spin}^T(n) \longrightarrow S^1 \times S^1$ is given by $l([g, z_1, z_2]) = (z_1^2, z_1 z_2)$.
- e. $p : \text{Spin}^T(n) \longrightarrow SO(n) \times S^1 \times S^1$ is given by $p([g, z_1, z_2]) = (\lambda(g), z_1^2, z_1 z_2)$. Hence, $p = \lambda^T \times l$. Here p is a 2-fold covering.

Thus, we obtain the following commutative diagram where the row and the column are exact.

$$\begin{array}{ccccccc}
 & & & & 1 & & \\
 & & & & \downarrow & & \\
 & & & & S^1 \times S^1 & & \\
 & & & & \downarrow j & \searrow & \\
 1 & \longrightarrow & \text{Spin}(n) & \xrightarrow{i} & \text{Spin}^T(n) & \xrightarrow{l} & S^1 \times S^1 \longrightarrow 1 \\
 & & \searrow \lambda & & \downarrow \lambda^T & & \\
 & & & & SO(n) & & \\
 & & & & \downarrow & & \\
 & & & & 1 & &
 \end{array}$$

Moreover, we have the following exact sequence:

$$1 \longrightarrow \mathbb{Z}_2 \longrightarrow \text{Spin}^T(n) \xrightarrow{p} SO(n) \times S^1 \times S^1 \longrightarrow 1.$$

Theorem 2 *The group $\text{Spin}^T(n)$ is isomorphic to $\text{Spin}^c(n) \times S^1$.*

Proof We define the map φ in the following way:

$$\begin{aligned} \text{Spin}(n) \times S^1 \times S^1 &\xrightarrow{\varphi} \text{Spin}^c(n) \times S^1 \\ (g, z_1, z_2) &\mapsto ([g, z_1], z_1 z_2) \end{aligned}$$

It can be easily shown that φ is a surjective homomorphism and the kernel of φ is $\{(1, 1, 1), (-1, -1, -1)\}$. Thus, the group $\text{Spin}^T(n)$ is isomorphic to $\text{Spin}^c(n) \times S^1$. \square

Since $\text{Spin}(n)$ is contained in the complex Clifford algebra $\mathbb{C}l_n$, the spin representation κ of the group $\text{Spin}(n)$ extends to a $\text{Spin}^T(n)$ -representation. For an element $[g, z_1, z_2]$ from $\text{Spin}^T(n)$ and any spinor $\psi \in \Delta_n$, the spinor representation κ^T of $\text{Spin}^T(n)$ is given by

$$\kappa^T[g, z_1, z_2]\psi = z_1^2 z_2 \kappa(g)(\psi).$$

Proposition 3 *If $n = 2k + 1$ is odd, then κ^T is irreducible.*

Proof Assume that $\{0\} \neq W \neq \Delta_{2k+1}$ is a Spin^T invariant subspace. Thus, we have $\kappa^T[g, z_1, z_2](W) \subseteq W$. That is, $z_1^2 z_2 \kappa(g)(W) \subseteq W$. In this case, for every $w \in W$ there exists a $w' \in W$ such that $z_1^2 z_2 \kappa(g)(w) = w'$. As $\kappa(g)(w) = \frac{1}{z_1^2 z_2} w' \in W$ and the representation κ of $\text{Spin}(n)$ is irreducible if n is odd, this is a contradiction. The representation κ^T of $\text{Spin}^T(n)$ has to be irreducible for $n = 2k + 1$. \square

Proposition 4 *If $n = 2k$ is even, then the spinor space Δ_{2k} decomposes into two subspaces $\Delta_{2k} = \Delta_{2k}^+ \oplus \Delta_{2k}^-$.*

Proof We know that the $\text{Spin}(n)$ representation Δ_{2k} decomposes into two subspaces Δ_{2k}^+ and Δ_{2k}^- . Thus, we obtain $z_1^2 z_2 \kappa(g)(\Delta_{2k}^+) \subseteq \Delta_{2k}^+$ and $z_1^2 z_2 \kappa(g)(\Delta_{2k}^-) \subseteq \Delta_{2k}^-$. Namely, $\kappa^T[g, z_1, z_2](\Delta_{2k}^+) \subseteq \Delta_{2k}^+$ and $\kappa^T[g, z_1, z_2](\Delta_{2k}^-) \subseteq \Delta_{2k}^-$. Hence, the $\text{Spin}^T(2k)$ representation Δ_{2k} decomposes into two subspaces Δ_{2k}^+ and Δ_{2k}^- . It can be easily seen that the $\text{Spin}^T(2k)$ representation Δ_{2k}^\pm is irreducible. \square

The Lie algebra of the group $\text{Spin}^T(n)$ is described by

$$\mathfrak{spin}^T(n) = \mathfrak{m}_2 \oplus i\mathbb{R} \oplus i\mathbb{R}.$$

The differential $p_* : \mathfrak{spin}^T(n) \rightarrow \mathfrak{so}(n) \oplus i\mathbb{R} \oplus i\mathbb{R}$ is defined by

$$p_*(e_\alpha e_\beta, \lambda i, \mu i) = (2E_{\alpha\beta}, 2\lambda i, (\lambda + \mu)i)$$

where λ and μ are any real numbers and $E_{\alpha\beta}$ is the $n \times n$ matrix with entries $(E_{\alpha\beta})_{\alpha\beta} = -1$, $(E_{\alpha\beta})_{\beta\alpha} = 1$ and all others are equal to zero. The inverse of the differential p_* is given by

$$p_*^{-1}(E_{\alpha\beta}, \lambda i, \mu i) = (\frac{1}{2}e_\alpha e_\beta, \frac{1}{2}\lambda i, (\mu - \frac{1}{2}\lambda)i).$$

3 Spin^T structure

Definition 5 A Spin^T structure on an oriented Riemannian manifold (M^n, g) is a $\text{Spin}^T(n)$ principal bundle $P_{\text{Spin}^T(n)}$ together with a smooth map $\Lambda : P_{\text{Spin}^T(n)} \rightarrow P_{SO(n)}$ such that the following diagram commutes:

$$\begin{array}{ccc} P_{\text{Spin}^T(n)} \times \text{Spin}^T(n) & \longrightarrow & P_{\text{Spin}^T(n)} \\ \downarrow \Lambda \times \lambda^T & & \downarrow \Lambda \\ P_{SO(n)} \times SO(n) & \longrightarrow & P_{SO(n)} \end{array}$$

From above definition we can construct a two-fold covering map

$$\Pi : P_{\text{Spin}^T(n)} \rightarrow P_{SO(n)} \times P_{S^1 \times S^1}.$$

Given a Spin^T structure $(P_{\text{Spin}^T(n)}, \Lambda)$, the map $\lambda^T : \text{Spin}^T(n) \rightarrow SO(n)$ induces an isomorphism

$$P_{\text{Spin}^T(n)} / S^1 \times S^1 \cong P_{SO(n)}.$$

In similar way, $\text{Spin}^T(n) / \text{Spin}(n) \cong S^1 \times S^1$ implies the isomorphism

$$P_{\text{Spin}^T(n)} / \text{Spin}(n) \cong P_{S^1 \times S^1}.$$

Note that on account of the inclusion map $i : \text{Spin}(n) \rightarrow \text{Spin}^T(n)$, every spin structure on M induces a Spin^T structure. Similarly, since there exists a inclusion map $\text{Spin}^c(n) \rightarrow \text{Spin}^T(n)$, every Spin^c structure on M induces a Spin^T structure.

4 Spinor bundle and Dirac operator

Let (M^n, g) be an oriented connected Riemannian manifold and $P_{SO(n)} \rightarrow M$ the $SO(n)$ -principal bundle of positively oriented orthonormal frames. The Levi-Civita connection ∇ on $P_{SO(n)}$ determine a connection 1-form ω on the principal bundle $P_{SO(n)}$ with values in $\mathfrak{so}(n)$, locally given by

$$\omega^e = \sum_{i < j} g(\nabla e_i, e_j) E_{ij}$$

where $e = \{e_1, \dots, e_n\}$ is a local section of $P_{SO(n)}$ and E_{ij} is the $n \times n$ matrix with entries $(E_{ij})_{ij} = -1$, $(E_{ij})_{ji} = 1$ and all others are equal to zero.

We fix a connection

$$(A, B) : TP_{S^1 \times S^1} \rightarrow i\mathbb{R} \oplus i\mathbb{R}$$

on the principal bundle $P_{S^1 \times S^1}$. The connections ω and (A, B) induce a connection

$$\omega \times (A, B) : T(P_{SO(n)} \times P_{S^1 \times S^1}) \rightarrow \mathfrak{so}(n) \oplus i\mathbb{R} \oplus i\mathbb{R}$$

on the fibre product bundle $P_{SO(n)} \times P_{S^1 \times S^1}$. Now we can define a connection 1-form $\omega \times \widetilde{(A, B)}$ on the principal bundle $P_{\text{Spin}^T(n)}$ such that the following diagram commutes:

$$\begin{array}{ccc} TP_{\text{Spin}^T(n)} & \xrightarrow{\omega \times \widetilde{(A, B)}} & \mathfrak{spin}^T(n) = \mathfrak{m}_2 \oplus i\mathbb{R} \oplus i\mathbb{R} \\ \downarrow \Pi_* & & \downarrow p_* \\ T(P_{SO(n)} \times P_{S^1 \times S^1}) & \xrightarrow{\omega \times (A, B)} & \mathfrak{so}(n) \oplus i\mathbb{R} \oplus i\mathbb{R} \end{array}$$

That is, the equality

$$p_* \circ \omega \times \widetilde{(A, B)} = \omega \times (A, B) \circ \Pi_*$$

holds.

Definition 6 *The spinor bundle of a Spin^T manifold is defined as the associated vector bundle*

$$\mathbb{S} = P_{\text{Spin}^T(n)} \times_{\kappa^T} \Delta_n$$

where $\kappa^T : \text{Spin}^T(n) \rightarrow GL(\Delta_n)$ is the spinor representation of $\text{Spin}^T(n)$. In case of $n = 2k$ the spinor bundle splits into the sum of two subbundles \mathbb{S}^+ and \mathbb{S}^- such that

$$\mathbb{S} = \mathbb{S}^+ \oplus \mathbb{S}^-, \quad \mathbb{S}^\pm = P_{\text{Spin}^T(n)} \times_{\kappa^{T\pm}} \Delta_n^\pm.$$

Any spinor field ψ can be identified with the map $\psi : P_{\text{Spin}^T(n)} \rightarrow \Delta_n$ satisfying the transformation rule $\psi(pg) = \kappa^T(g^{-1})\psi(p)$. The absolute differential of a section ψ with respect to $\omega \times \widetilde{(A, B)}$ determines a covariant derivative

$$\tilde{\nabla} : \Gamma(\mathbb{S}) \rightarrow \Gamma(T^*M \otimes \mathbb{S})$$

given by

$$\tilde{\nabla}\psi = d\psi + \kappa_{*1}^T(\omega \times \widetilde{(A, B)})\psi$$

where $\kappa_{*1}^T : \mathfrak{spin}^T(n) \rightarrow \text{End}(\Delta_n)$ is the derivative of κ at the identity $1 \in \text{Spin}^T(n)$. It can be also shown that

$$\kappa_{*1}^T(e_\alpha e_\beta, \lambda i, \mu i) = \kappa(e_\alpha e_\beta) + (2\lambda i + \mu i)Id$$

where λ and μ are any real numbers and κ is the spin representation of the group $\text{Spin}(n)$.

Now we give the local formulas for connections. Fix a section $s : U \rightarrow P_{S^1 \times S^1}$ of the principal bundle $P_{S^1 \times S^1}$. Then, we obtain the local connection form

$$(A^s, B^s) : TU \rightarrow i\mathbb{R} \oplus i\mathbb{R}$$

where $A^s, B^s : TU \rightarrow i\mathbb{R}$. $e \times s : U \rightarrow P_{SO(n)} \times P_{S^1 \times S^1}$ is a local section of the fiber product bundle $P_{SO(n)} \times P_{S^1 \times S^1}$. $\widetilde{e \times s}$ is a lift of this section to the

two-fold covering $\Pi : P_{Spin^T(n)} \rightarrow P_{SO(n)} \times P_{S^1 \times S^1}$. The local connection form $\omega \times \widetilde{(A, B)}^{(\widetilde{e \times s})}$ on the principal bundle $P_{Spin^T(n)}$ is given by the formula

$$\omega \times \widetilde{(A, B)}^{(\widetilde{e \times s})} = \left(\frac{1}{2} \sum_{i < j} g(\nabla e_i, e_j) e_i e_j, \frac{1}{2} A^s, B^s - \frac{1}{2} A^s \right)$$

Hence, this connection form induces a connection $\widetilde{\nabla}$ on the spinor bundle \mathbb{S} . We can locally describe $\widetilde{\nabla}$ by

$$\widetilde{\nabla}_X \psi = d\psi(X) + \frac{1}{2} \sum_{i < j} g(\nabla_X e_i, e_j) e_i e_j \psi + \frac{1}{2} A^s \psi + B^s \psi \quad (1)$$

where $\psi : U \rightarrow \Delta_n$ is a section of the spinor bundle \mathbb{S} .

Definition 7 *The first order differential operator*

$$D = \mu \circ \widetilde{\nabla} : \Gamma(\mathbb{S}) \xrightarrow{\widetilde{\nabla}} \Gamma(T^*M \otimes \mathbb{S}) \xrightarrow{\mu} \Gamma(\mathbb{S})$$

where μ denotes Clifford multiplication, is called the Dirac operator.

The Dirac operator D is locally given by

$$D\psi = \sum_{i=1}^n e_i \cdot \widetilde{\nabla}_{e_i} \psi \quad (2)$$

where $\{e_1, \dots, e_n\}$ is a local orthonormal frame on the manifold M .

The Dirac operator has the following property:

Theorem 8 *Let f be a smooth function and $\psi \in \Gamma(\mathbb{S})$ be a spinor field. Then,*

$$D(f \cdot \psi) = (grad f \cdot \psi) + f D\psi.$$

Proof By using the definition of the Dirac operator D we can compute $D(f \cdot \psi)$ as follows:

$$\begin{aligned} D(f \cdot \psi) &= \sum_{i=1}^n e_i \cdot \widetilde{\nabla}_{e_i} (f \cdot \psi) \\ &= \sum_{i=1}^n e_i \cdot (e_i(f) \cdot \psi + f \widetilde{\nabla}_{e_i} \psi) \\ &= \sum_{i=1}^n e_i(f) e_i \cdot \psi + f \sum_{i=1}^n e_i \cdot \widetilde{\nabla}_{e_i} \psi \\ &= (grad f) \cdot \psi + f D\psi \end{aligned}$$

□

Now we can define the Laplace operator on the spinor bundle \mathbb{S} .

Definition 9 *Let $\psi \in \Gamma(\mathbb{S})$ be a spinor field. The Laplace operator Δ on spinors is defined by*

$$\Delta\psi = - \sum_{i=1}^n \left(\widetilde{\nabla}_{e_i} \widetilde{\nabla}_{e_i} \psi + div(e_i) \widetilde{\nabla}_{e_i} \psi \right). \quad (3)$$

4.1 Schrödinger-Lichnerowicz type formula

The square D^2 of the Dirac operator and the Laplace operator Δ are second order differential operators. We derive Schrödinger-Lichnerowicz type formula by computing their difference $D^2 - \Delta$.

The curvature $R^{\mathbb{S}}$ of the spinor covariant derivative $\tilde{\nabla}$ is an $End(\mathbb{S})$ valued 2-form by

$$R^{\mathbb{S}}(X, Y)\psi = \tilde{\nabla}_X \tilde{\nabla}_Y \psi - \tilde{\nabla}_Y \tilde{\nabla}_X \psi - \tilde{\nabla}_{[X, Y]}\psi$$

where $\psi \in \Gamma(\mathbb{S})$ and $X, Y \in \Gamma(TM)$. Now we want to describe $R^{\mathbb{S}}$ in terms of the curvature tensor R .

Let $\Omega^\omega : TP_{SO(n)} \times TP_{SO(n)} \rightarrow \mathfrak{so}(n)$ be the curvature form of the Levi-Civita connection with the components

$$\Omega^\omega = \sum_{i < j} \Omega_{ij} E_{ij}$$

where $\Omega_{ij} : TP_{SO(n)} \times TP_{SO(n)} \rightarrow \mathbb{R}$. The commutative diagram defining the connection $\omega \times (A, B)$ implies that the curvature form of $\omega \times (A, B)$ is

$$\Omega^{\omega \times (A, B)} = \frac{1}{2} \sum_{i < j} \Pi^*(\Omega_{ij}) e_i e_j \oplus \frac{1}{2} \Pi^*(dA) \oplus \Pi^*(dB).$$

Hence the 2-form $R^{\mathbb{S}}$ with values in the spinor bundle \mathbb{S} is obtained by the following formula:

$$R^{\mathbb{S}}(., .)\psi = \frac{1}{2} \sum_{i < j} \Omega_{ij} e_i e_j \cdot \psi + \frac{1}{2} dA \cdot \psi + dB \cdot \psi.$$

Let $\{e_1, \dots, e_n\}$ be orthonormal frame field, $\Omega_{ij}(X, Y) = g(R(X, Y)e_i, e_j)$ the components of the curvature form of the Levi-Civita connection,

$X = \sum_{k=1}^n X^k e_k$ and $Y = \sum_{l=1}^n Y^l e_l$ be vector fields on the Riemannian manifold M . Then we have

$$\begin{aligned} \Omega_{ij}(X, Y) &= g(R(X, Y)e_i, e_j) \\ &= \sum_{k, l=1}^n R_{kl ij} X^k Y^l \\ &= \sum_{k, l=1}^n R_{kl ij} e^k(X) e^l(Y) \\ &= \frac{1}{2} \sum_{k, l=1}^n R_{kl ij} (e^k \wedge e^l)(X, Y). \end{aligned}$$

where $\{e^1, \dots, e^n\}$ is the frame dual to $\{e_1, \dots, e_n\}$. Thus, we obtain the following local formula for the curvature form

$$\Omega^{\omega \times (A, B)} = \frac{1}{4} \sum_{i < j} \sum_{k, l=1}^n R_{kl ij} (e^k \wedge e^l) e_i e_j + \frac{1}{2} dA + dB$$

and the 2-form $R^S(.,.)$ is calculated as follows:

$$R^S(.,.)\psi = \frac{1}{4} \sum_{i < j} \sum_{k,l=1}^n R_{klij}(e^k \wedge e^l) e_i e_j \cdot \psi + \frac{1}{2} dA \cdot \psi + dB \cdot \psi.$$

By using the above properties of the curvature form R^S on spinor bundle S we deduce the following result:

Proposition 10 *Let Ric be the Ricci tensor. Then, the following relation holds:*

$$\sum_{\alpha=1}^n e_\alpha \cdot R^S(X, e_\alpha) \psi = -\frac{1}{2} Ric(X) \cdot \psi + \frac{1}{2} (X \lrcorner dA) \cdot \psi + (X \lrcorner dB) \cdot \psi \quad (4)$$

Proof In [2] it is proved the following relation:

$$\sum_{\alpha=1}^n \sum_{i < j} \sum_{k,l=1}^n R_{klij}(e^k \wedge e^l) e_\alpha e_i e_j \cdot \psi = -2 Ric(X) \cdot \psi \quad (5)$$

It can be easily seen the following two relations:

$$\sum_{\alpha=1}^n e_\alpha \cdot dA(X, e_\alpha) \cdot \psi = (X \lrcorner dA) \cdot \psi \quad (6)$$

and

$$\sum_{\alpha=1}^n e_\alpha \cdot dB(X, e_\alpha) \cdot \psi = (X \lrcorner dB) \cdot \psi. \quad (7)$$

Then, using (5), (6) and (7), we obtain the claimed equivalence. \square

Now, we derive Schrödinger-Lichnerowicz-type formula in the following way:

Proposition 11 *Let s be scalar curvature of the Riemannian manifold and let $dA = \Omega^A$ and $dB = \Omega^B$ be the imaginary-valued 2-forms of the connections (A, B) in the $(S^1 \times S^1)$ -bundle associated with $Spin^T$ structure. Then, we have the following formula:*

$$D^2\psi = \Delta\psi + \frac{s}{4}\psi + \frac{1}{2}dA \cdot \psi + dB \cdot \psi.$$

Proof

$$\begin{aligned} D^2\psi &= \sum_{i,j} e_i \cdot \tilde{\nabla}_{e_i}(e_j \cdot \tilde{\nabla}_{e_j}\psi) \\ &= \sum_{i,j} e_i \cdot \nabla_{e_i} e_j \cdot \tilde{\nabla}_{e_j}\psi + e_i e_j \cdot \tilde{\nabla}_{e_i} \tilde{\nabla}_{e_j}\psi \\ &= \sum_{i,j,k} g(\nabla_{e_i} e_j, e_k) e_i e_k \cdot \tilde{\nabla}_{e_j}\psi + \sum_{i,j} e_i e_j \cdot \tilde{\nabla}_{e_i} \tilde{\nabla}_{e_j}\psi \\ &= \Delta\psi + \sum_{j, i \neq k} g(\nabla_{e_i} e_j, e_k) e_i e_k \cdot \tilde{\nabla}_{e_j}\psi + \sum_{i \neq j} e_i e_j \cdot \tilde{\nabla}_{e_i} \tilde{\nabla}_{e_j}\psi \end{aligned} \quad (8)$$

Now we can calculate the following sum:

$$\begin{aligned}
 \sum_{i \neq k} g(\nabla_{e_i} e_j, e_k) e_i e_k &= - \sum_{i \neq k} g(e_j, \nabla_{e_i} e_k) e_i e_k \\
 &= - \sum_{i < k} g(e_j, \nabla_{e_i} e_k - \nabla_{e_k} e_i) e_i e_k \\
 &= \sum_{i < k} g(e_j, [e_k, e_i]) e_i e_k
 \end{aligned}$$

From (8) we get

$$\begin{aligned}
 D^2 \psi &= \Delta \psi + \sum_{j, i < k} g(e_j, [e_k, e_i]) e_i e_k \tilde{\nabla}_{e_j} \psi + \sum_{i < j} e_i e_j \cdot (\tilde{\nabla}_{e_i} \tilde{\nabla}_{e_j} \psi - \tilde{\nabla}_{e_j} \tilde{\nabla}_{e_i} \psi) \\
 &= \Delta \psi + \sum_{i < j} e_i e_j (\tilde{\nabla}_{e_i} \tilde{\nabla}_{e_j} \psi - \tilde{\nabla}_{e_j} \tilde{\nabla}_{e_i} \psi - \tilde{\nabla}_{[e_i, e_j]} \psi) \\
 &= \Delta \psi + \frac{1}{2} \sum_{i, j} e_i e_j R^S(e_i, e_j) \psi.
 \end{aligned}$$

Using the identity (4) and multiplying by e_i we deduce

$$\begin{aligned}
 D^2 \psi &= \Delta \psi - \frac{1}{4} \sum_i e_i \text{Ric}(e_i) \cdot \psi + \frac{1}{4} \sum_i e_i \cdot (e_i \lrcorner dA) \cdot \psi + \frac{1}{2} \sum_i e_i \cdot (e_i \lrcorner dB) \cdot \psi \\
 &= \Delta \psi + \frac{s}{4} \psi + \frac{1}{2} dA \cdot \psi + dB \cdot \psi.
 \end{aligned}$$

□

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